

# Inorganic Analogues of Cyclic Organic Molecules

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Received March 26, 1974/July 1, 1974

*Ab initio* SCFMO calculations are given for inorganic systems related to the hydrocarbon rings  $C_3H_3^+$ ,  $C_3H_6$ ,  $C_4H_4$ ,  $C_4H_8$ , and  $C_6H_6$ . The molecules studied belong to the same symmetry group and/or are isoelectric with the hydrocarbon rings, and include  $H_3^+$ ,  $H_4$ ,  $H_6$ ,  $N_3^+$ ,  $N_4$ ,  $N_6$ ,  $O_3$ ,  $O_4$ , and several others. Preferred geometries and energies are reported for single- $\zeta$  and double- $\zeta$  basis sets. The relationship between the hydrocarbon rings and their inorganic analogues is shown. Although lone-pair repulsions destabilize the inorganic rings, their possible existence as molecules, intermediates, and transition states is discussed.

*Key word:* Inorganic ring systems with equal atoms

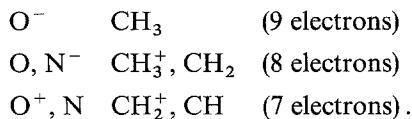
## 1. Introduction

The predictive power of Walsh's rules [1, 2] for determining molecular geometry is well-established for triatomic systems of the form  $AH_2$ ,  $AB_2$ , and  $HAB$ . It has been shown by Buenker and Peyerimhoff [3] that it is possible to extend Walsh's rules to more general systems of substituted triatomics, such as  $H_nAB_2$ , in order to make geometry predictions for more complex molecules. This is true because hydrogen substitution can alter the position of valence orbital energies without altering their angular dependence, and the preferred bond angle is determined by following the variation of sums of orbital energies. Thus, for example, it was shown that the energy variation on angle bending in ozone is quite similar to that in the isoelectronic molecule  $C_3H_6$ , with minima occurring in both cases near  $60^\circ$  (ring) and  $120^\circ$  (bent) geometries [4–7].

In this paper we consider the possible inorganic analogues of the hydrocarbon rings  $C_3H_3^+$  (cyclopropenyl cation),  $C_3H_6$  (cyclopropane),  $C_4H_4$  (cyclobutadiene),  $C_4H_8$  (cyclobutane), and  $C_6H_6$  (benzene). These molecules have received considerable attention because, for example,  $C_3H_3^+$  and  $C_6H_6$  form prototype aromatic molecules while  $C_4H_4$  is antiaromatic [8];  $C_3H_6$  and  $C_4H_8$  are also important systems in organic chemistry and have been thoroughly studied by molecular orbital techniques [9–11]. We shall make use of the well known properties of such molecules in the present study.

We can classify organic and inorganic systems as analogues provided that

- 1) They are isoelectronic. The following replacements can be considered:



- 2) They belong to the same symmetry group (isosymmetric).

When these two conditions are fulfilled, it is highly likely that the organic and inorganic molecules will have the same electronic configuration, and hence a similar geometry (some exceptions will be discussed). Under certain circumstances, which have been discussed in the recent review by Buenker and Peyerimhoff [45], the isosymmetric requirement is unnecessary. This allows extension of results obtained for ABA systems, for example, to include ABC systems.

To illustrate the use of the two criteria of equivalence given above, consider a possible inorganic analogue of cyclopropane ( $D_{3h}$  symmetry, 24 electrons) which is cyclic ozone ( $D_{3h}$ , 24 electrons). Figure 1 shows a system of localized orbitals where the equivalence between CH bonds and oxygen lone pairs is revealed; the high symmetry ensures a comparable orbital energy diagram for both molecules.

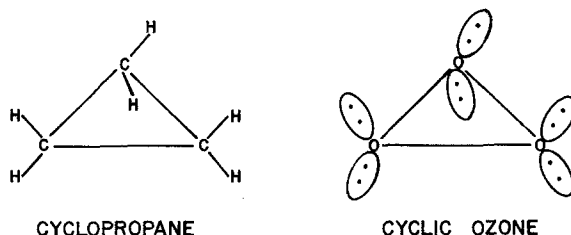
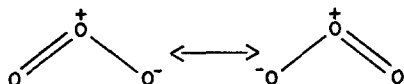
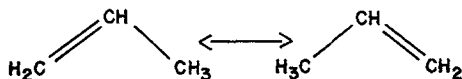


Fig. 1. Localized orbitals ( $sp^3$  hybridization assumed) for cyclopropane and cyclic ozone

The point which emerges from the analogy is that ozone should exist as a *stable ring*. In fact, the ground state geometry of ozone is bent at  $116^\circ$  [12], but there is a deep secondary minimum for cyclic ozone at  $60^\circ$  [4-7]. Inspection of resonance structures for bent ozone (below),



and simple replacement by isoelectronic groups ( $O$ ,  $CH_2$ ), ( $O^+$ ,  $CH$ ), ( $O^-$ ,  $CH_3$ ) leads to a prediction in the reverse direction – there should be a stable hydrocarbon of form



which, of course, is propylene, a structural isomer of cyclopropane<sup>1</sup>. Note that in this case we have dropped requirement [2] of symmetry equivalence and still retained some predictive power. In other cases we can drop the requirement that the systems be isoelectronic; as will be shown for the hydrogen rings the requirement of symmetry equivalence (belonging to the same point group) can lead to interesting analogies with the hydrocarbon rings. We now begin a systematic investigation of the inorganic analogues.

<sup>1</sup> There are some traps to be avoided in this process. Ozone has equal bond lengths, propylene does not. Propylene cannot be represented by resonance structures since the presence of hydrogens makes the positions non-equivalent. Another error can occur where there are partially occupied degenerate orbitals; compare  $O_2$  and ethylene.

## 2. Method of Calculation

Calculations reported in this paper use either a minimum basis (single- $\zeta$ , hereafter called SZ) or double- $\zeta$  (DZ) basis of Slatertype exponential orbitals in an *ab initio* SCFMO calculation. The molecular orbital program used was obtained from Stevens and has been discussed in recent publications [13, 14] (the integral package is obtainable from the Quantum Chemistry Program Exchange as QCPE 161).

Table 1. Computed total energies for the cyclic molecules and ions

Molecule	$r$ (bohr)	Basis set	Energy (hartree)	Comment	Ref.
H <sub>2</sub>	1.40	SZ	- 1.1281	$\zeta = 1.20$	<sup>a</sup>
	1.40	DZ	- 1.1283	$\zeta = 0.90, 1.20$	<sup>a</sup>
	1.40	1s, 1s', 2p	- 1.1669	full CI	[25]
	1.40		- 1.1745	exact	[20]
H <sub>3</sub> <sup>+</sup>	1.66	SZ <sup>e</sup>	- 1.2699	$\zeta = 1.40$	<sup>a</sup>
	1.66	large	- 1.3397	full CI	[19]
	1.66	large	- 1.342	est. ( $\pm 0.001$ )	[19]
H <sub>4</sub>	2.30	SZ <sup>e</sup>	- 1.9834	<sup>3</sup> A <sub>2g</sub> , $\zeta = 1.029$	<sup>a</sup>
	2.30	DZ	- 2.0079	<sup>3</sup> A <sub>2g</sub> , 0.90, 1.20	<sup>a</sup>
	2.40	1s, 1s', 2p	- 2.0909	<sup>3</sup> A <sub>2g</sub> , full CI	[25]
H <sub>6</sub>	2.06	SZ <sup>e</sup>	- 3.2081	$\zeta = 1.10$	<sup>a</sup>
	1.86 <sup>e</sup>	DZ	- 3.2408	$\zeta = 0.90, 1.20$	[26]
N <sub>2</sub>	2.178 <sup>e</sup>	SZ	-108.5694	Clementi ex.	[27]
	2.075 <sup>e</sup>	DZ	-108.8644	Huzinaga ex.	[27]
	2.0675	large	-108.9928	near H-F	<sup>b</sup>
N <sub>3</sub> <sup>+</sup>	2.580 <sup>e</sup>	SZ	-162.389	ion-optimized	[27]
	2.655 <sup>e</sup>	DZ	-162.597	Huzinaga ex.	[27]
N <sub>4</sub>	2.714 <sup>e</sup>	SZ	-216.925	<sup>3</sup> A <sub>2g</sub> , Clem. ex.	[27]
	2.700	DZ	-217.240	<sup>3</sup> A <sub>2g</sub> , Huzin. ex.	[27]
N <sub>6</sub>	2.606 <sup>e</sup>	SZ	-325.468	Clementi ex.	[27]
O <sub>2</sub>	2.280	SZ	-149.0906	<sup>3</sup> $\Sigma_g^-$ , Clem. Ex.	<sup>a</sup>
	2.280	SZ	-149.0825	<sup>3</sup> $\Sigma_g^-$ , 7.66, 2.28	<sup>a</sup>
	2.280	DZ	-149.5623	<sup>3</sup> $\Sigma_g^-$ , Clementi ex.	<sup>a</sup>
	2.280	large	-149.6639		<sup>c</sup>
O <sub>3</sub>	2.670 <sup>e</sup>	SZ	-223.5568		[5]
	2.730	DZ	-224.1793		[7]
	2.730	DZ <sup>+</sup>	-224.2794		[7]
O <sub>4</sub>	2.67 <sup>e</sup>	SZ	-298.1081	$\zeta = 7.66, 2.28$	<sup>a</sup>
	2.84 <sup>e</sup>	DZ	-298.9546	Clementi ex.	<sup>a</sup>
C <sub>3</sub> H <sub>6</sub>	2.76	GTO	-116.8974		<sup>d</sup>
C <sub>4</sub> H <sub>4</sub>	2.691	GTO	-153.42	<sup>3</sup> A <sub>2g</sub>	[41]
C <sub>4</sub> H <sub>8</sub>	2.929	SZ	-155.8370		[11]

<sup>a</sup> This work.

<sup>b</sup> Cade, P. E., Sales, K. D., Wahl, A. C.: J. Chem. Phys. **44**, 1973 (1966).

<sup>c</sup> Schaefer, H. F.: J. Chem. Phys. **54**, 2207 (1971).

<sup>d</sup> Siu, A. K. Q., St. John 3rd, W. M., Hayes, E. F.: J. Am. Chem. Soc. **92**, 7249 (1970).

<sup>e</sup> Optimized.

For neutral molecules we use Clementi's [15] best-atom exponents for SZ basis sets and Huzinaga [16] best-atom exponents for DZ basis sets, unless otherwise noted. Some optimization was done to find the best exponents for ions. Where open-shell situations are encountered, two techniques are used: Nesbet's method of symmetry and equivalence restrictions [17] in an open shell calculation, or a limited configuration interaction. All of the computed energies, together with the results of other workers, have been collected in Table 1.

### 3. Hydrogen Rings

Consider the symmetric rings of hydrogen atoms having molecular formulas  $H_3^+$ ,  $H_4$ , and  $H_6$ . Respective symmetries for the rings are  $D_{3h}$ ,  $D_{4h}$ , and  $D_{6h}$  so by criterion [2] we look for analogies with the isosymmetric hydrocarbon rings  $C_3H_3^+$ ,  $C_4H_4$ , and  $C_6H_6$ . In Fig. 2 the orbital energy patterns (based on SZ results), which are completely symmetry determined, are shown. These orbital energies were obtained for the geometries and exponents listed in Table 1, which are nearly optimum. Inspection of the orbital energies with reference to the energy of a hydrogen 1s orbital ( $-0.5$  hartree) shows that, from orbital energy considerations alone,  $H_2$ ,  $H_3^+$ , and  $H_6$  should be stable while  $H_4$  is unstable. Exactly the same energy level diagram forms the basis for the Hückel  $4n + 2$  rule for the stability of  $pi$ -systems, from which follows the stability of  $C_3H_3^+$  and  $C_6H_6$  (aromatic), and instability of  $C_4H_4$  (antiaromatic). We now ask: Is there evidence, experimental or theoretical, concerning the stability of hydrogen rings?

The  $H_3^+$  ion has been observed in the mass spectrometer [18]. From data on appearance potentials, the estimated enthalpy change for  $H_2 + H^+ \rightarrow H_3^+$  is  $-115$  kcal/mole. In an extensive quantum mechanical calculation of the

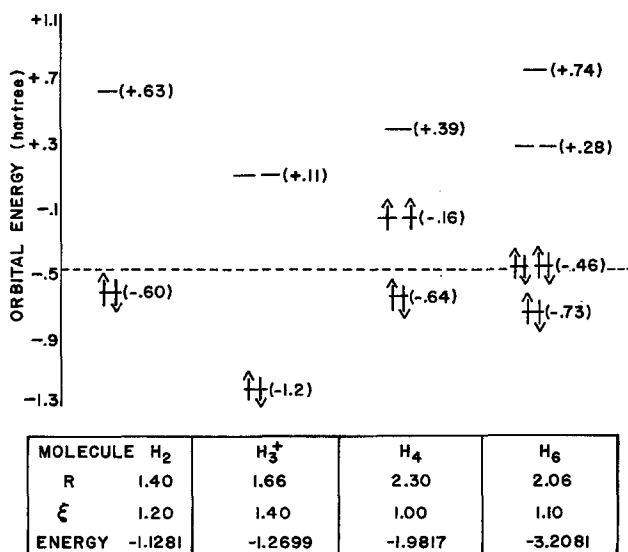


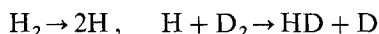
Fig. 2. SZ orbital energies for hydrogen rings. Dashed line = energy of a hydrogen 1s orbital ( $-0.5$  hartree)

preferred geometry of  $H_3^+$  by Csizmadia *et al.* [19], by extrapolation of their CI result they obtained

$$\text{Energy of } H_3^+ (\text{exact}) = -1.342 \pm 0.001 \text{ hartree.}$$

The preferred geometry was cyclic with  $r = 1.66$  bohr. Since the exact energy of  $H_2$  [20] is  $-1.1745$  hartree (at the potential minimum), the computed energy difference from  $H_2 + H^+$  is 105 kcal/mole, in reasonable agreement with the experimental value of 115. Thus, the  $H_3^+$  ion forms a stable ring.

The search for a low energy  $H_4$  molecule has provided a challenge to theoreticians, following the shock tube study by Bauer and coworkers [21] of the  $H_2 + D_2 \rightarrow 2HD$  isotopic exchange reaction. They observed the low activation energy of 42 kcal/mole [21, 22] and proposed a *molecular* exchange mechanism, involving a vibrationally excited  $H_2(D_2)$  in a four-center transition state. The atomic exchange mechanism,



was considered to be ruled out on kinetic grounds. An exhaustive search by *ab initio* techniques [23–25] for any four-center transition state which lies below the energy required for the atomic exchange mechanism has never revealed such a transition state. It now appears unlikely that such a transition state will be found.

The reason for the instability of square planar  $H_4$  is evident from inspection of Fig. 2, where the “antiaromatic” nature of the occupied  $H_4$  orbitals is shown. By contrast,  $H_6$  appears stable and this led the author to propose the possibility that the exchange reaction proceeds through a *six-center* transition state [26]. DZ calculations on the energy of  $H_2$ ,  $H_4$ , and  $H_6$  showed  $E(H_4) - 2E(H_2) = 156$  kcal/mole, i.e.  $H_4$  is a high energy form, but  $E(H_6) - 3E(H_2) = 90$  kcal/mole, which is below the dissociation energy of  $H_2$ . The reaction was studied along both concerted termolecular and two-step bimolecular paths. In each case the relatively low energy of  $H_6$  was the essential part of the mechanism. Thus, the hydrogen rings show useful analogies to the hydrocarbon rings, a fact which follows from their symmetry-determined orbital energy patterns. Although the true mechanism of this reaction must still be considered to be in doubt, some recent molecular beam results [49] on reactions of this type lend support to the idea of a 6-center transition state.

#### 4. Nitrogen Rings

Using a system of hybrid orbitals common to both nitrogen and carbon, the nitrogen rings and their organic analogues are shown in Fig. 3. Because the two systems are isoelectronic and isosymmetric, their MO energy level diagrams will be quite similar. Where the form of the MO's is completely symmetry determined, as is the case for the *pi*-MO's, the pattern must be identical. By analogy, then,  $N_3^+$  and  $N_6$  should form stable aromatic systems whereas  $N_4$ , like  $C_4H_4$ , should be antiaromatic and unstable. However,  $N_2$  is the stable form of nitrogen and obviously a closer examination is in order.

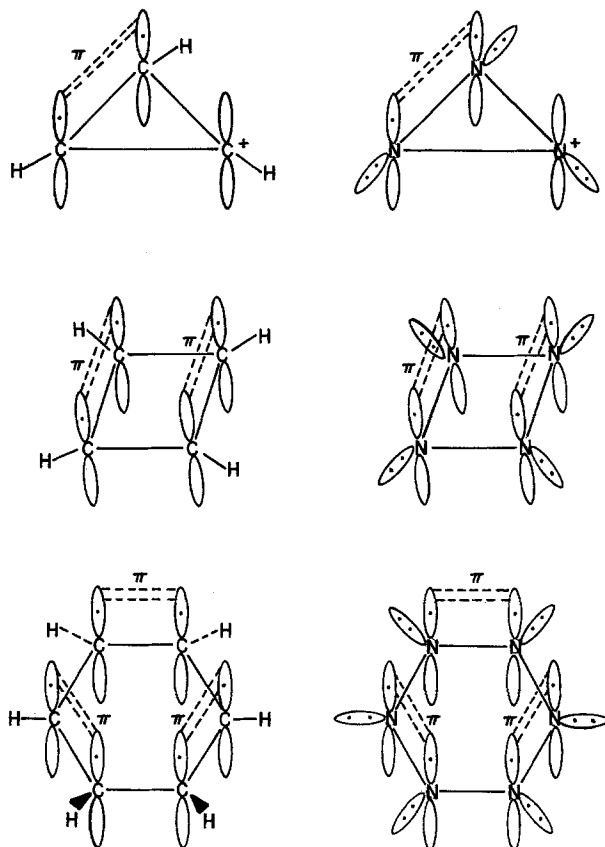
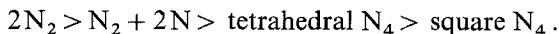


Fig. 3.  $N_3^+$ ,  $N_4$ , and  $N_6$  and their organic analogues cyclopropenyl $^+$ , cyclobutadiene, and benzene

An extensive study of these rings [27] led to the following conclusions (energies and optimized distances appear in Table 1):

1) Although the most stable form of  $N_3^+$  is linear and asymmetric [28], there is a secondary minimum for cyclic  $N_3^+$ .

2)  $N_4$  is highly destabilized relative to two molecules of  $N_2$ . In a DZ basis, the observed order of stability is



Relative energies in hartree are  $2N_2$  (0.0); tetrahedral (0.436), square planar triplet (0.488). The  $N_4$  species lie above the dissociation energy of  $N_2$  (0.360 hartree).

3)  $N_6$  lies 0.240 hartree above  $3N_2$  in a SZ basis and shows some stabilization relative to  $N_4$  in the same basis.

4) SZ basis sets appear to seriously overestimate the stability of the small rings, when compared to DZ calculations. For example, for  $N_4$  in an SZ basis relative energies are  $2N_2$  (0.0), rectangular singlet (0.204), square planar triplet (0.214), tetrahedral (0.218).

These results can be understood by considering the nature of the bonds  $N\equiv N$ ,  $N=N$ , and  $N-N$ . The extremely stable triple bond, and rather unstable double and single bonds have been explained [29, 30] on the basis of lone pair repulsions, e.g. as in diimide (I) and hydrazine (II).



In fact the ground state geometry of I is trans [46] and of II is skewed [47]; for these geometries the lone pair repulsions are clearly lessened but apparently still quite destabilizing (thus leading to weakened double and single bonds). In this regard, it is striking that protonation of  $N_2H_4$  to produce  $N_2H_5^+$ , followed by another protonation to give the  $N_2H_6^{++}$  ion leads to a series of bond distances which are 1.47, 1.45, and 1.42 Å respectively [43], apparently due to the lessening of lone-pair repulsive effects on protonation. This effect of the sigma-system apparently swamps the more modest energy effect coming from the stability, or lack of it, in the *pi*-system. Nevertheless, the secondary minimum for cyclic  $N_3^+$  and the relative stability of  $N_6$  compared to  $N_4$  can be understood by reference to the hydrocarbon analogues.

Since the cyclobutadiene molecule is expected to be more stable in its rectangular singlet rather than square triplet form, we might look for a similar stabilization in rectangular  $N_4$ . However, the SZ calculations show a stabilization of only 0.01 hartree relative to the square triplet. Thus a DZ calculation on rectangular  $N_4$  would be expected to give an energy ca. 0.48 hartree above  $N_2 + N_2$ . The lack of CH bonding, or, alternatively, the presence of interacting lone pairs in the sigma-system, leads to a significant destabilization in cyclic  $N_4$  relative to cyclobutadiene [27, 41].

It was predicted by Longuet-Higgins and Orgel [31] that cyclobutadiene would be stabilized by interaction with a transition metal, primarily through a stabilizing interaction with the degenerate *pi* non-bonding orbital on  $C_4H_4$  and the metal *d*-orbitals. The subsequent synthesis of cyclobutadieneiron tricarbonyl by Pettit and coworkers [32] verified this hypothesis. It has been proposed by Shusturovich *et al.* [33, 34] that a complex of the form

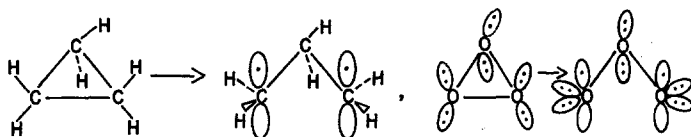


may have exceptional stability, and could be involved in the fixation of molecular nitrogen. Since the *pi*-orbitals are essentially identical with those in cyclobutadiene (although slightly lower in energy), the argument is well-founded. However, the results given in this paper on  $N_4$  render it unlikely that such a complex could achieve sufficient stability through this interaction to overcome the inherent instability of the  $N_4$  ring.

## 5. Oxygen Rings

The relationship between ozone and cyclopropane was discussed by Buenker and Peyerimhoff [3]. The existence of a double minimum in the bending potential for ozone has been shown by several authors [4–7], where one form corresponds to the cyclic ( $60^\circ$ ) molecule and a second form to the bent ( $\sim 116^\circ$ ) geometry. Again, the SZ calculations drastically overemphasize the ring stability [5] but the DZ calculations predict an energy almost competitive with that of bent ozone [4, 6, 7]. The chemistry of ozone has been discussed on this basis [5, 6].

The cyclization of ozone involves an orbital crossing from an occupied  $\pi$  non-bonding orbital to an occupied sigma bonding orbital. This can be seen, in reverse direction, to be analogous to the ring opening of cyclopropane to produce a trimethylene diradical (below):



Although bent ozone is stable and trimethylene is not, there should be an activation energy for each process. SZ calculations by Salem and coworkers [9, 10] on cyclopropane for the process indicated give a barrier height of 0.083 hartree, while for cyclic  $\rightarrow$  bent ozone the barrier height is 0.134 hartree in a SZ calculation [5], 0.050 hartree in a DZ calculation [4]. The analogy is useful for understanding the electronic changes in the reaction, as shown in more detail by Buenker and Peyerimhoff [3, 48].

It may be noted here that optimized SZ calculations [5] predict cyclic ozone to be 48 kcal/mole *more* stable than bent ozone, while in a CNDO/2 calculation [7] cyclic ozone is preferred by 300 kcal/mole! A GVB calculation [6] favors the bent species by approximately 35 kcal/mole, whereas a DZ calculation favors bent ozone by ca. 10 kcal/mole [4, 7]. It is probable that calculations beyond the Hartree-Fock limit would further favor the bent form, since in that geometry CI will be more important due to the presence of low lying vacant orbitals [5, 44].

The  $O_4$  molecule is isoelectronic with cyclobutane and may therefore possess an energy minimum in the square planar form. All experimental measurements on  $O_4$  [35–37] suggest that it is a loosely bound dimer of  $O_2$ , consisting of two slightly perturbed  $O_2$  molecules of ill-defined geometry. More to the present interest is the observation of the isotopic exchange reaction  $^{16}O_2 + ^{18}O_2 \rightarrow 2^{16}O-^{18}O$  in shock tube studies [38, 39], the mechanism of which may be either atomic, via dissociation of  $O_2$  and then  $^{16}O + ^{18}O_2 \rightarrow ^{16}O-^{18}O + ^{18}O$ , or molecular, via a four center transition state. The experimental data show that a molecular mechanism is preferred [38]. We report here calculations on the stability of the four membered ring  $O_4$ , to test the feasibility of the exchange pathway through this structure.

Results for a square planar ring are given in Table 1. Once again the SZ calculations favor the ring relative to the DZ calculations, and by a substantial



Table 2. Molecular orbitals of cyclobutane and O<sub>4</sub> (D<sub>4h</sub> symmetry)

Orbital symmetry	Orbital type	C <sub>4</sub> H <sub>8</sub> orbital energy (eV) <sup>a</sup>	O <sub>4</sub> orbital energy (eV) <sup>b</sup>
1a <sub>1g</sub>	1s	-304.850	-562.068
1e <sub>u</sub>	1s	-304.849	-562.076
1b <sub>2g</sub>	1s	-304.847	-562.079
2a <sub>1g</sub>	σ	-29.539	-44.387
2e <sub>u</sub>	σ	-23.678	-35.305
2b <sub>2g</sub>	σ	-19.504	-25.835
1a <sub>2u</sub>	π <sub>b</sub>	-17.287	-18.688
3a <sub>1g</sub>	σ	-16.712	-17.202
1e <sub>g</sub>	π <sub>nb</sub>	-14.133	-13.507
1b <sub>1g</sub>	σ	-12.715	-17.893
1b <sub>1u</sub>	π <sub>ab</sub>	-11.624	-8.941
3e <sub>u</sub>	σ	-11.050	-12.677
2a <sub>2u</sub>	π*	+13.972	none
4e <sub>u</sub>	σ*	+15.194	+9.390
4a <sub>1g</sub>	σ*	+16.699	none
1a <sub>2g</sub>	σ*	+17.036	+13.175
2e <sub>g</sub>	π*	+18.895	none
5e <sub>u</sub>	σ*	+19.289	none
3b <sub>2g</sub>	σ*	+19.926	+13.469
2b <sub>1u</sub>	π*	+23.567	none
4b <sub>2g</sub>	σ*	+24.005	none

<sup>a</sup> Reference [11].

<sup>b</sup> 1s = 7.66, 2s = 2p = 2.28, r = 2.67 bohr.

amount. Ignoring the question of correlation energy problems due to the open shell in O<sub>2</sub>, simple subtraction in the same basis set gives; for 2O<sub>2</sub> → O<sub>4</sub>,

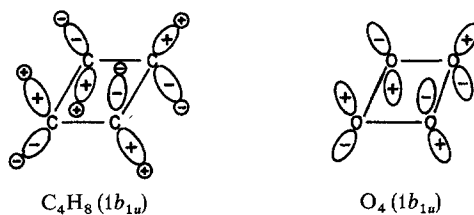
$$\Delta E(\text{SZ}) = -298.1081 - 2(-149.0825) = 0.0567 \text{ hartree} = 36 \text{ kcal/mole},$$

$$\Delta E(\text{DZ}) = -298.9546 - 2(-149.5623) = 0.170 \text{ hartree} = 107 \text{ kcal/mole}.$$

Both calculations show an O<sub>4</sub> energy below that of the experimental O<sub>2</sub> dissociation energy, 118 kcal/mole [40]. It was found that distortion of the O<sub>4</sub> molecule along the normal modes of the square [11] leads to an increase in the total energy along each mode, so the square form represents a potential minimum, not a transition state (SZ basis). The magnitude of the barrier to be overcome during the decomposition of cyclic O<sub>4</sub>, and hence its computed stability, will be the subject of a future publication (which will include a search for the lowest energy exchange pathway).

The extent of the analogy between O<sub>4</sub> and C<sub>4</sub>H<sub>8</sub> can be seen by inspection of the orbital energies obtained in a SZ calculation for each. These orbitals and their symmetry designations in D<sub>4h</sub> [11] are shown in Table 2. The oxygen orbitals are of course shifted to lower energy. The ordering is generally parallel, except that the π<sub>nb</sub> (pi non-bonding, not shown) and π<sub>ab</sub> (pi antibonding, below) orbitals in

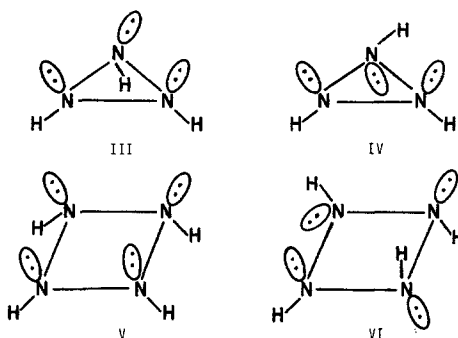
$O_4$  are destabilized relative to the corresponding orbitals in  $C_4H_8$  ( $1e_g$  and  $1b_{1u}$  orbitals, Table 2). This same effect



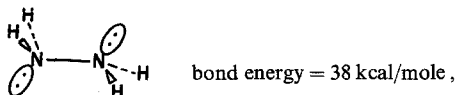
was observed in a study of  $C_4H_4$  [41], where it was shown that the H-atoms play an important role in stabilizing the  $\sigma_{cc}$  antibonding molecular orbitals. No such possibility exists in  $O_4$ , and the  $1b_{1u}$  MO, at  $-8.94$  eV, lies well above the rest.

## 6. Other Rings

These arguments can be readily extended to mixed rings and larger systems. For example, the hypothetical molecules  $N_3H_3$  and  $N_4H_4$  are isoelectronic with  $C_3H_6$  and  $C_4H_8$  and may be expected to show minima for comparable geometries (below):

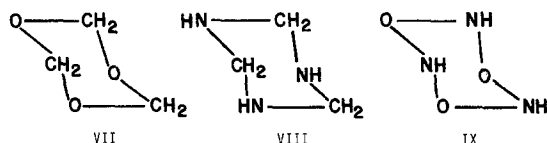


Structures III–V exhibit strong lone pair repulsions but VI should be relatively favored. However, the very weak single bond in hydrazine, which can also adopt a comparable geometry (below), renders all of these structures unstable. Why hydrazine,



and also diimide, both of which can minimize lone pair repulsions by adoption of the trans structure, should be so unstable remains an interesting unsolved problem. The simple picture of lone pair-repulsions is insufficient, and one would like to know the extent of the lone pair-bond pair interaction.

A mixed 6-membered ring, trioxymethylene (VII, below) is isoelectronic with cyclohexane and probably exists in the chair conformation also. This ring forms in a low activation energy



concerted polymerization involving three monomers of formaldehyde [42]. An analogous possibility for a nitrogen substituted system suggests itself (VIII) or even a mixed HNO system (IX), but in these cases the monomers  $\text{H}_2\text{C}=\text{NH}$  and  $\text{HN}=\text{O}$  are unstable and the reactions have not been observed.

## 7. Summary and Conclusions

By analogy with hydrocarbon rings through the establishment of a) isoelectronic, and b) isosymmetric criteria, one can discuss the stability of some hypothetical, but potentially observable molecules. Thus  $\text{O}_3$  exists as a secondary minimum in cyclic form and  $\text{O}_4$  could be involved in isotopic exchange; cyclic  $\text{N}_3^+$  and  $\text{N}_6$  show signs of stability whereas  $\text{N}_4$  is very unstable;  $\text{H}_3^+$  and  $\text{H}_6$  are favored while  $\text{H}_4$  is not, and extensions may be made to larger systems. The effect of lone pair repulsions reduces the stability of nitrogen and oxygen rings, which may nevertheless exist as minima or transition states in the reactions discussed. The presence of NH and OH bonds will serve to reduce these high energies by stabilizing the framework antibonding orbitals [41], and protonation of lone pairs can also add a stabilizing effect.

*Acknowledgement.* The author would like to thank the National Research Council of Canada (Grant No. C-1013) for financial support.

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